

A Frequency-Domain Condition for the Stability of Feedback Systems Containing a Single Time-Varying Nonlinear Element

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It is proved that a condition similar to the Nyquist criterion guarantees the stability (in an important sense) of a large class of feedback systems containing a single time-varying nonlinear element. In the case of principal interest, the condition is satisfied if the locus of a certain complex-valued function (a) is bounded away from a particular disk located in the complex plane, and (b) does not encircle the disk.

I. INTRODUCTION

The now well-known techniques introduced by Lyapunov have led to many very interesting results concerning the stability of time-varying nonlinear feedback systems governed by systems of differential equations. However, these methods have by no means led to a definitive theory of stability for even the simplest nontrivial time-varying nonlinear feedback systems. The general problem is, of course, one of considerable difficulty.

The unparalleled utility of the Nyquist stability criterion for single-loop, linear, time-invariant feedback systems is directly attributable to the fact that it is an explicit frequency-domain condition. The Nyquist locus not only indicates the stability or instability of a system, it presents the information in such a way as to aid the designer in arriving at a suitable design. The criterion is useful even in cases in which the system is so complicated that a sufficiently accurate analysis is not feasible, since experimental measurements can be used to construct the loop-gain locus.

The primary purpose of this article is to point out that some recently

obtained mathematical results,¹ not involving the theory of Lyapunov, imply that a condition similar to, and possessing the advantages of, the Nyquist criterion guarantees the stability (in an important sense) of feedback systems containing a single time-varying nonlinear element.*†

II. THE PHYSICAL SYSTEM AND DEFINITION OF \mathcal{L}_2 -STABILITY

Consider the feedback system of Fig. 1. We shall restrict our discussion throughout to cases in which g_1 , f , u , and v denote real-valued measurable functions of t defined for $t \geq 0$.

The block labeled ψ is assumed to represent a memoryless time-varying (not necessarily linear) element that introduces the constraint $u(t) = \psi[f(t), t]$, in which $\psi(x, t)$ is a function of x and t with the

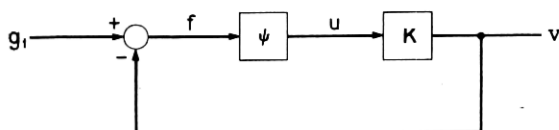


Fig. 1 — Nonlinear feedback system.

properties that $\psi(0, t) = 0$ for $t \geq 0$ and there exist a positive constant β and a real constant α such that

$$\alpha \leq \frac{\psi(x, t)}{x} \leq \beta, \quad t \geq 0$$

for all real $x \neq 0$. In particular, we permit the extreme cases in which $\psi(x, t)$ is either independent of t or linear in x [i.e., $\psi(x, t) = \psi(1, t)x$].

The block labeled K represents the linear time-invariant portion of the forward path. It is assumed to introduce the constraint

$$v(t) = \int_0^t k(t - \tau)u(\tau)d\tau - g_2(t), \quad t \geq 0$$

in which k and g_2 are real-valued functions such that

$$\int_0^\infty |k(t)| dt < \infty, \quad \int_0^\infty |g_2(t)|^2 dt < \infty. \quad (1)$$

* The results of Ref. 1 relate to feedback systems containing an arbitrary finite number of time-varying nonlinear elements, but, with the exception of the case discussed here, they do not admit of a simple geometric interpretation.

† For results concerned with frequency-domain conditions for the global asymptotic stability (a sense of stability that is different from the one considered here) of nonlinear systems, see, for example, Refs. 2-4.

The function g_2 takes into account the initial conditions at $t = 0$. Our assumptions regarding \mathbf{K} are satisfied, for example, if, as is often the case, u and v are related by a differential equation of the form

$$\sum_{n=0}^N a_n \frac{d^n v}{dt^n} = \sum_{n=0}^{N-1} b_n \frac{d^n u}{dt^n}, \quad t \geq 0$$

in which the a_n and the b_n are constants with $a_N \neq 0$, and

$$\sum_{n=0}^N a_n s^n \neq 0 \quad \text{for} \quad \operatorname{Re}[s] \geq 0.$$

However, we *do not* require that u and v be related by a differential equation (or by a system of differential equations).

Assumption: We shall assume throughout that the response v is well defined and satisfies the inequality

$$\int_0^t |v(\tau)|^2 d\tau < \infty \quad (2)$$

for all finite $t > 0$, for each initial-condition function g_2 that meets the conditions stated above and each input g_1 such that

$$\int_0^\infty |g_1(t)|^2 dt < \infty.$$

Although this assumption plays an important role in the proof of the theorem to be presented, from an engineering viewpoint it is a trivial restriction (see Ref. 5).

Definition: We shall say that the feedback system of Fig. 1 is " \mathcal{L}_2 -stable" if and only if there exists a positive constant ρ with the property that the response v satisfies

$$\left(\int_0^\infty |v(t)|^2 dt \right)^{\frac{1}{2}} \leq \rho \left(\int_0^\infty |g_1(t) + g_2(t)|^2 dt \right)^{\frac{1}{2}} + \left(\int_0^\infty |g_2(t)|^2 dt \right)^{\frac{1}{2}}$$

for every initial-condition function g_2 that meets the conditions stated above, and every input g_1 such that

$$\int_0^\infty |g_1(t)|^2 dt < \infty.$$

In particular, if the system is \mathcal{L}_2 -stable, then the response is square-integrable whenever the input is square-integrable.

It can be shown* that the response $v(t)$ approaches zero as $t \rightarrow \infty$ for any square-integrable input g_1 , provided that the system is \mathcal{L}_2 -stable,

* See the proof of Theorem 6 of Ref. 1.

$g_2(t) \rightarrow 0$ as $t \rightarrow \infty$, and

$$\int_0^{\infty} |k(t)|^2 dt < \infty. \quad (3)$$

In addition, it follows at once from the Schwarz inequality that the response $v(t)$ is uniformly bounded on $[0, \infty)$ for any square-integrable input g_1 , provided that the system is \mathcal{L}_2 -stable, $g_2(t)$ is uniformly bounded on $[0, \infty)$, and (3) is satisfied.

III. SUFFICIENT CONDITIONS FOR \mathcal{L}_2 -STABILITY

Theorem: Let

$$K(i\omega) = \int_0^{\infty} k(t) e^{-i\omega t} dt, \quad -\infty < \omega < \infty.$$

The feedback system of Fig. 1 is \mathcal{L}_2 -stable if one of the following three conditions is satisfied:

(i) $\alpha > 0$; and the locus of $K(i\omega)$ for $-\infty < \omega < \infty$ (a) lies outside the circle C_1 of radius $\frac{1}{2}(\alpha^{-1} - \beta^{-1})$ centered on the real axis of the complex plane at $[-\frac{1}{2}(\alpha^{-1} + \beta^{-1}), 0]$, and (b) does not encircle C_1 (see Fig. 2)

(ii) $\alpha = 0$, and $\text{Re}[K(i\omega)] > -\beta^{-1}$ for all real ω

(iii) $\alpha < 0$, and the locus of $K(i\omega)$ for $-\infty < \omega < \infty$ is contained within the circle C_2 of radius $\frac{1}{2}(\beta^{-1} - \alpha^{-1})$ centered on the real axis of the complex plane at $[-\frac{1}{2}(\alpha^{-1} + \beta^{-1}), 0]$ (see Fig. 3).

Proof: Note first that

$$\int_0^{\infty} |u(t)|^2 dt \leq \max(\beta^2, |\alpha|^2) \int_0^{\infty} |f(t)|^2 dt,$$

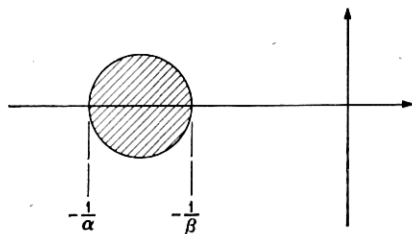


Fig. 2 — Location of the “critical circle” C_1 in the complex plane ($\alpha > 0$). The feedback system is \mathcal{L}_2 -stable if the locus of $K(i\omega)$ for $-\infty < \omega < \infty$ lies outside C_1 and does not encircle C_1 .

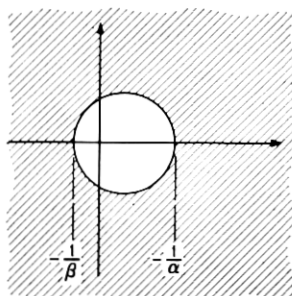


Fig. 3 — Location of the “critical circle” C_2 in the complex plane ($\alpha < 0$). The feedback system is \mathfrak{L}_2 -stable if the locus of $K(i\omega)$ for $-\infty < \omega < \infty$ is contained within C_2 .

and hence, by a well-known result,

$$\begin{aligned} \int_0^\infty \left| \int_0^t k(t-\tau)u(\tau)d\tau \right|^2 dt &\leq \left(\int_0^\infty |k(t)| dt \right)^2 \int_0^\infty |u(t)|^2 dt \\ &\leq \max(\beta^2, |\alpha|^2) \left(\int_0^\infty |k(t)| dt \right)^2 \int_0^\infty |f(t)|^2 dt. \end{aligned}$$

Using Minkowski's inequality,

$$\begin{aligned} \left(\int_0^\infty |v(t)|^2 dt \right)^{\frac{1}{2}} &\leq \left(\int_0^\infty \left| \int_0^t k(t-\tau)u(\tau)d\tau \right|^2 dt \right)^{\frac{1}{2}} \\ &\quad + \left(\int_0^\infty |g_2(t)|^2 dt \right)^{\frac{1}{2}} \leq \max(\beta, |\alpha|) \int_0^\infty |k(t)| dt \\ &\quad \cdot \left(\int_0^\infty |f(t)|^2 dt \right)^{\frac{1}{2}} + \left(\int_0^\infty |g_2(t)|^2 dt \right)^{\frac{1}{2}}. \end{aligned}$$

Consider now the relation between $(g_1 + g_2)$ and f :

$$g_1(t) + g_2(t) = f(t) + \int_0^t k(t-\tau)\psi[f(\tau),\tau]d\tau, \quad t \geq 0$$

and suppose that

$$\int_0^\infty |g_1(t) + g_2(t)|^2 dt < \infty.$$

According to the results of Ref. 1, our assumptions* imply that there

* In Ref. 1 it is assumed that

$$\int_0^t |f(\tau)|^2 d\tau < \infty$$

exists a positive constant ρ_1 (which does not depend upon g_1 or g_2) such that

$$\int_0^\infty |f(t)|^2 dt < \rho_1 \int_0^\infty |g_1(t) + g_2(t)|^2 dt$$

provided that, with

$$K(s) = \int_0^\infty k(t)e^{-st} dt$$

and $\omega = \text{Im}[s]$,

- (i) $1 + \frac{1}{2}(\alpha + \beta)K(s) \neq 0$ for $\text{Re}[s] \geq 0$, and
- (ii) $\frac{1}{2}(\beta - \alpha) \max_{-\infty < \omega < \infty} |K(i\omega)[1 + \frac{1}{2}(\alpha + \beta)K(i\omega)]^{-1}| < 1$.

Thus the feedback system of Fig. 1 is \mathcal{L}_2 -stable if conditions (i) and (ii) are satisfied.

According to the well-known theorem of complex-function theory that leads to the Nyquist criterion, condition (i) is satisfied if (and only if) the polar plot of $K(i\omega)$ for $-\infty < \omega < \infty$ does not encircle or pass through the point $[-2(\alpha + \beta)^{-1}, 0]$. It can easily be verified that condition (ii) is met if one of the following three conditions is satisfied.

(a) $\alpha > 0$, and the locus of $K(i\omega)$ for $-\infty < \omega < \infty$ lies outside the circle C_1 of radius $\frac{1}{2}(\alpha^{-1} - \beta^{-1})$ centered in the complex plane at $[-\frac{1}{2}(\alpha^{-1} + \beta^{-1}), 0]$.

(b) $\alpha = 0$, and $\text{Re}[K(i\omega)] > -\beta^{-1}$ for all real ω .

(c) $\alpha < 0$, and the locus of $K(i\omega)$ for $-\infty < \omega < \infty$ is contained within the circle C_2 of radius $\frac{1}{2}(\beta^{-1} - \alpha^{-1})$ centered in the complex plane at $[-\frac{1}{2}(\alpha^{-1} + \beta^{-1}), 0]$.

If $\alpha > 0$, the point $[-2(\alpha + \beta)^{-1}, 0]$ lies on the real-axis diameter of C_1 , while if condition (b) or (c) is met, it is impossible for the polar plot of $K(i\omega)$ to encircle the point $[-2(\alpha + \beta)^{-1}, 0]$. Therefore, the conditions of the theorem guarantee that the feedback system is \mathcal{L}_2 -stable.

Remarks

With regard to the necessity of our sufficient conditions for \mathcal{L}_2 -stability, consider, for example, the case in which $\alpha > 0$ and suppose, for simplicity, that v and u are related by a differential equation of the type mentioned in Section II. Then, a moment's reflection shows that there exists a $\psi(x, t)$, in fact a $\psi(x, t)$ which is independent of t and linear in x ,

for all finite $t > 0$. Our assumption that (2) is satisfied for all finite $t > 0$ implies that this condition is met.

that satisfies our assumptions and for which the feedback system is *not* \mathcal{L}_2 -stable, provided that for some value of ω , $K(i\omega)$ is a point on the real-axis diameter of C_1 . This clearly shows that the condition is in the correct "ball park." Similar remarks can be made concerning our conditions for the cases in which $\alpha < 0$ and $\alpha = 0$.

IV. FURTHER PROPERTIES OF THE FEEDBACK SYSTEM OF FIG. 1

It is possible to say much more about the properties of the feedback system on the basis of frequency-domain information if our assumptions regarding $\psi(x, t)$ are strengthened.

For example, suppose that

$$\alpha \leq \frac{\psi(x_1, t) - \psi(x_2, t)}{x_1 - x_2} \leq \beta, \quad \psi(0, t) = 0 \quad (4)$$

for $t \geq 0$ and all real $x_1 \neq x_2$, and that one of the three conditions of our theorem is met. Let g_1 and \hat{g}_1 denote two arbitrary input functions such that

$$\int_0^t |g_1(\tau)|^2 d\tau < \infty \quad \text{and} \quad \int_0^t |\hat{g}_1(\tau)|^2 d\tau < \infty$$

for all finite $t > 0$, and

$$\int_0^\infty |g_1(\tau) - \hat{g}_1(\tau)|^2 d\tau < \infty.$$

Let v and \hat{v} , respectively, denote the (assumed well defined) responses due to g_1 and \hat{g}_1 . Then if

$$\int_0^t |v(\tau)|^2 d\tau < \infty \quad \text{and} \quad \int_0^t |\hat{v}(\tau)|^2 d\tau < \infty$$

for all finite $t > 0$, and the assumptions of Section II are met, it follows* that

$$\int_0^\infty |v(\tau) - \hat{v}(\tau)|^2 d\tau < \infty$$

and that there exists a positive constant λ (which does not depend upon g_1 or \hat{g}_1) such that

$$\int_0^\infty |v(\tau) - \hat{v}(\tau)|^2 d\tau \leq \lambda \int_0^\infty |g_1(\tau) - \hat{g}_1(\tau)|^2 d\tau.$$

* Consider Theorem 1 of Ref. 6 with $h_1(t) = f_1(t) = 0$ for $t < 0$.

Suppose now that $\psi(x, t)$ satisfies (4) and is either independent of t or periodic in t with period T for each x , and that one of the three conditions of our theorem is met. Assume that the initial-condition function $g_2(t)$ approaches zero as $t \rightarrow \infty$, and that the input $g_1(t)$ applied at $t = 0$ is a bounded periodic function with period T . Then it can be shown* that there exists a bounded periodic function p , with period T , which is independent of g_2 and such that the (assumed well defined) response $v(t)$ approaches $p(t)$ as $t \rightarrow \infty$, provided that the conditions of Section II are met, (2) is satisfied for all finite $t > 0$, and

$$\int_0^\infty \left| \int_t^\infty |k(\tau)| d\tau \right|^2 dt < \infty, \quad \int_0^\infty |(1+t)k(t)|^2 dt < \infty. \quad (5)$$

Observe that the conditions of (5) are satisfied if u and v are related by a differential equation of the form described in Section II.

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* See Theorem 3 of Ref. 6.